

phosphate (KDP) crystal (see Fig. 1) [7]. The photons are separated using a polarizing BS and coupled into single-mode fibers. Detection of one photon is used to herald its twin, which is sent to a 50:50 BS to produce the single-photon path-entangled reference state. To create a time-dependent controllable phase, one path contains a piezoelectric translator (PZT) which changes the relative lengths of the two paths. The PZT is driven with a triangle wave at a frequency of 600 Hz and has an amplitude corresponding to 4π of phase. A square pulse is sent to a time-to-digital converter (TDC) at the start of each period. Each path is guided to separate fiber beam splitters (FBS).

To simulate a star in the laboratory, part of the pump laser pulse is propagated with a known transverse intensity profile to a rotating diffuser followed by a linear polarizer to create a thermal-like state ($g^{(2)}(0)=1.93\pm 0.03$) with a near-Lambertian angular scattering distribution. The diffuse light propagates for 1 m to a 50:50 BS, with each output monitored with a bare, single-mode fiber that forms the telescope. One fiber is translated across the wavefront, sampling over a 6.2-mm range in 0.01-mm steps, while the other fiber remains at a fixed position. Each fiber (T1 and T2 in Fig. 1) receives an average of 0.0038 photons per pulse and connects to the remaining input port of one FBS. The outputs of the FBSs and the fiber collecting the herald photon are monitored with superconducting nanowire single-photon detectors and a TDC. To ensure temporal overlap, Hong-Ou-Mandel interference is found by monitoring the coincidences of each FBS while changing the relative delays of each input [8].

The collected coincidences are filtered so that only coincidences (these can be either heralded or unheralded) between different paths are considered. The probability of such a coincidence is given by

$$P(|v|, \delta, \phi) = (1 \pm |v| \cos(\delta - \phi))/2, \quad (1)$$

where $|v|$ and ϕ are the magnitude and phase of the complex visibility at a given baseline respectively, and phase gained from transmission or reflection at the beam splitters determines the sign. Finally, the phase (δ_j) applied for each coincidence count can be measured given the relative timing to the PZT position. The phasor ($e^{i\delta_j}$) is averaged over many events to retrieve the visibility [9].

3. Data & Results

A double slit source was created using two 0.5 mm-wide slits separated by 1 mm. For each baseline position, the absolute value of the visibility was measured ten times. The average of each measurement is shown in Fig. 2a) for unheralded and heralded data and their fit to theory. Note that the visibility is much higher for the heralded case where the single-photon nature of the local oscillator does not include shot noise. A reconstruction of the source intensity distribution autocorrelation via Fourier transform of the visibility squared is shown in Fig. 2b). Here the separation of the side peaks from the main peak shows the separation of the slits, while the half width of the side peak shows the width of an individual slit. In Fig. 2c), the signal-to-noise ratio per coincidence event is shown for each baseline value, confirming that more information was gathered per event in the heralded case compared to unheralded. In summary, we have shown reconstruction of the source intensity distribution autocorrelation, a strict advantage of a non-local oscillator over a classical local oscillator scheme, and an important application of a future quantum network allowing for longer baselines in visible wavelength VLBI.

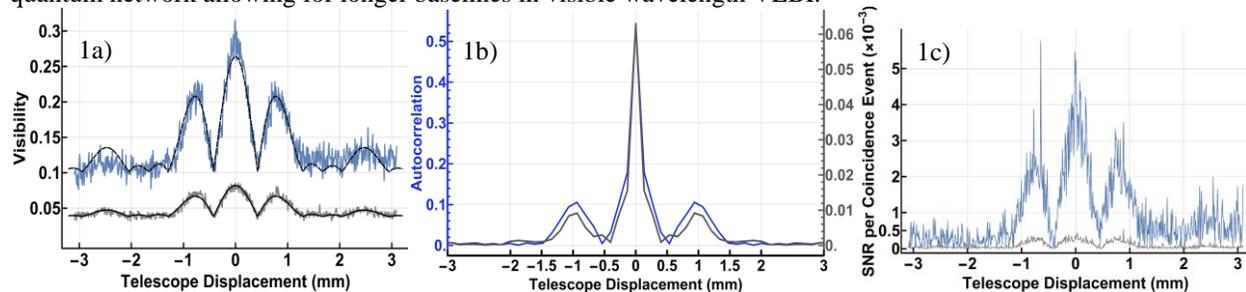


Fig. 2: a) visibility vs baseline, b) source distribution autocorrelation, and c) SNR per coincidence event. Blue(grey) is (un)heralded; black represents fitted theory curves.

- [1] Thalacker, C., et al. "Anonymous and secret communication in quantum networks." *New J. Phys.* 23, 083026 (2021)
- [2] Broadbent, A; Fitzsimons, J; Kashefi, E. "Universal Blind Quantum Computation." 50th Annual IEEE Symposium on Foundations of Computer Science, 517–526 (2009).
- [3] Kómár, P. et al. "A quantum network of clocks." *Nature Physics.* 10, 582–587 (2014).
- [4] Gottesman, D.; Jennewein, T.; and Croke, S., "Longer-baseline telescopes using quantum repeaters." *Phys. Rev. Lett.* 109, 070503 (2012)
- [5] Khabiboulline, E. T. et al, "Quantum-assisted telescope arrays." *Phys. Rev. A* 100, 022316 (2019).
- [6] Tsang, M. "Quantum nonlocality in weak-thermal-light interferometry." *Phys. Rev. Lett.* 107(27), 270402 (2011).
- [7] Mosely, P. J., et al. "Heralded Generation of Ultrafast Single Photons in Pure Quantum States." *Phys. Rev. Lett.* 100, 133601 (2008)
- [8] C. K. Hong, Z. Y. Ou, and L. Mandel, "Measurement of subpicosecond time intervals between two photons by interference," *Phys. Rev. Lett.* 59, 2044–2046 (1987).
- [9] Tamimi, A., et al., "Fluorescence-detected Fourier transform electronic spectroscopy by phase-tagged photon counting." *Opt. Exp.* 28, 25194–25214 (2020).